

Typicality in quantum statistical mechanics

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Open-Minded

Content

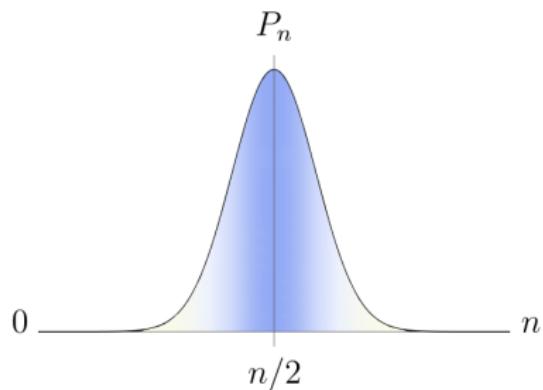
Introduction - What is Typicality?

Dynamical Quantum Typicality

Transition probabilities

Conclusion

Typicality



Gedankenexperiment:
Tossing a (fair) coin n times

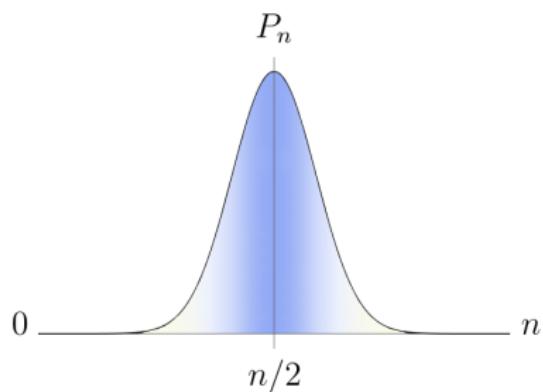
Consider the following result:

$$\{1, 1, \dots\} \in \{0, 1\}^n$$

n small (e.g. $n = 2$): no surprise

n large (e.g. $n = 10^{23}$): highly unlikely or rather impossible

Typicality



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- ▶ The boundary between typical and untypical results is blurred.
- ▶ Prerequisite: A large sample space (here large n)

Approach to statistical physics

Usual approach

Macrostate ρ determined by several fixed macroscopic variables (energy, magnetization, ...) and maximization of entropy

time evolution

$$\partial_t \rho = -i[H, \rho]$$

expectation values

$$\langle A \rangle = \text{Tr } \rho A$$

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Alternative approach

Microstates in **well defined regions** show very similar properties

Choosing a typical $|\psi\rangle$ will be sufficient

time evolution

$$i\partial_t |\psi\rangle = H|\psi\rangle$$

expectation values

$$\langle A \rangle = \langle \psi | A | \psi \rangle$$

Dynamical Quantum Typicality

Finite dimensional Hilbert space $\mathcal{H} = \mathbb{C}^N$

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$$M_a = \{|\varphi\rangle \in \mathcal{H} \mid \langle\varphi|\varphi\rangle = 1 \wedge \langle\varphi|A|\varphi\rangle \approx a\} \text{ then}$$

$$\forall t : |\varphi\rangle, |\varphi'\rangle \in M_a : \langle\varphi|B(t)|\varphi\rangle \approx \langle\varphi'|B(t)|\varphi'\rangle$$

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- ▶ This statement is not dependent on the details of the dynamics.
- ▶ Requirements: A large Hilbert space dimension N and ...

Hilbert space averaging

Goal: Find the typical $f(\psi)$ (e.g. expectation value) for random quantum states $|\psi\rangle$

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Hilbert space average

$$\text{HA}[f(\psi)] := \frac{\int d[\psi] f(\psi) \delta(\langle\psi|\psi\rangle - 1)}{\int d[\psi] \delta(\langle\psi|\psi\rangle - 1)}$$

Confinement to unit sphere by normalization $\langle\psi|\psi\rangle = 1$

Expectation values

To be more precise

Sampling $|\psi\rangle$ from a **uniform distribution** resp. integration over
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 $\psi_i \rightarrow -\psi_i \Rightarrow \forall i \neq j : \text{HA} [\psi_i^* \psi_j] = 0$

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 $\psi_i \rightarrow -\psi_i \Rightarrow \forall i \neq j : \text{HA} [\psi_i^* \psi_j] = 0$
- ▶ $\text{HA} [\sum_i |\psi_i|^2] = 1 \Rightarrow \forall i : \text{HA} [|\psi_i|^2] = \frac{1}{N}$

$$\Rightarrow \text{HA} [\langle \psi | A | \psi \rangle] = \frac{\text{Tr } A}{N}$$

Changing the distribution

Up to now uniform distribution

Density matrix ρ defines a new set of states $|\phi\rangle := \sqrt{N\rho} |\psi\rangle$

Ref.:

C. Bartsch and J. Gemmer: Phys. Rev. Lett. **102** 110403 (2009)

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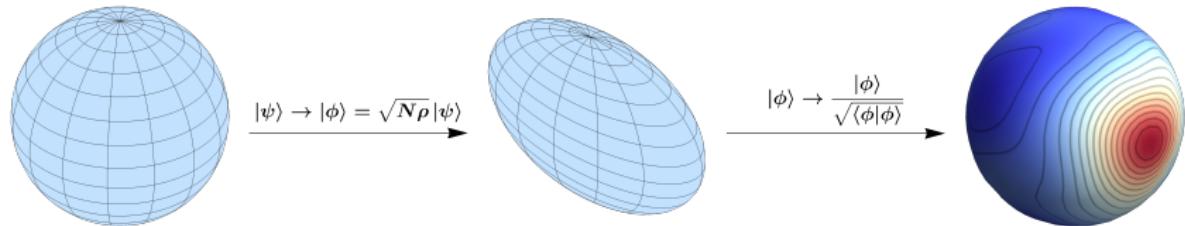
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Geometrically speaking



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Validity of Typicality

Hilbert space variance

$$\begin{aligned}\text{HV} [\langle \psi | A | \psi \rangle] &:= \text{HA} [\langle \psi | A | \psi \rangle^2] - \text{HA} [\langle \psi | A | \psi \rangle]^2 \\ &= \frac{1}{N+1} \left(\frac{\text{Tr } A^2}{N} - \frac{(\text{Tr } A)^2}{N^2} \right)\end{aligned}$$

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For $|\phi\rangle = \sqrt{N\rho} |\psi\rangle$ and **Purity** $P = \text{Tr } \rho^2$

$$\text{HV} [\langle \phi | \phi \rangle] \leq P \quad \text{HV} [\langle \phi | A | \phi \rangle] \leq \|A\|^2 P$$

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Dynamical Quantum Typicality

Choose ρ_a so that $\text{Tr } \rho_a A = a$

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$$\text{HV}[\langle\phi|A|\phi\rangle] \leq \|A\|^2 P \quad \text{HV}[\langle\phi|B(t)|\phi\rangle] \leq \|B\|^2 P$$

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Retrieve Dynamical Quantum Typicality for low Purity

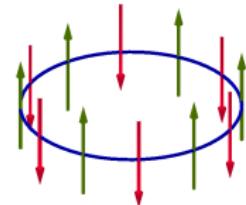
⇒ When is $P \ll 1$?

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Example: Kicked Ising Chain



n -particle spin-1/2-chain ($N = 2^n$)

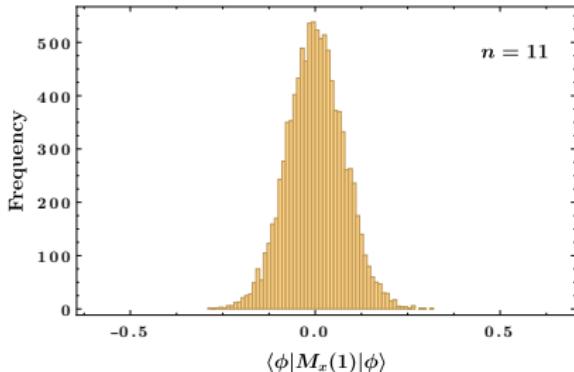
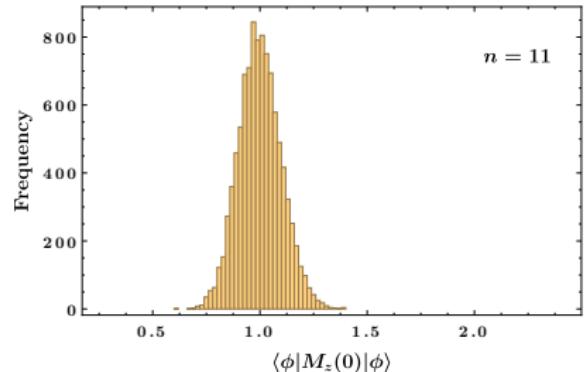
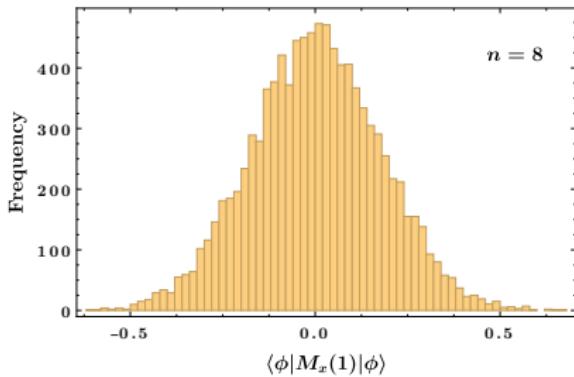
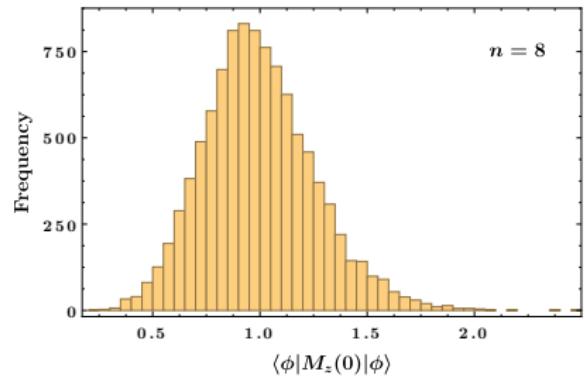
$$H = \sum_{i=1}^n J \sigma_i^z \sigma_{i+1}^z + h \sigma_i^z + \sum_{i=1}^n b \sigma_i^x \sum_{\tau=-\infty}^{\infty} \delta(t - \tau)$$

with **discrete time evolution**

$$U = U_I U_K = e^{-iH_I} e^{-iH_K}$$

Magnetization $M_\alpha = \sum_{i=1}^n \sigma_i^\alpha$ $\langle M_\alpha \rangle \in [-n, n]$

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Transition probabilities

Now $f(\chi, \psi) = |\langle \chi | A | \psi \rangle|^2$ and $|\chi\rangle, |\psi\rangle$ uniformly distributed, A unitary

$$\text{HA}_{\psi, \chi} \left[|\langle \chi | A | \psi \rangle|^2 \right] = \text{HA}_\chi \left[\frac{\langle \chi | AA^\dagger | \chi \rangle}{N} \right] = \frac{\text{Tr } AA^\dagger}{N^2} = \frac{1}{N}$$

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Impose additional condition: $\langle \chi | \psi \rangle = z$ with $|z| \leq 1$

First leave $|\chi\rangle$ fixed

$$\text{HA}_\psi \left[|\langle \chi | A | \psi \rangle|^2 \delta(\langle \chi | \psi \rangle - z) \right] = \frac{\int d[\psi] |\langle \chi | A | \psi \rangle|^2 \delta(\langle \chi | \psi \rangle - z) \delta(\langle \psi | \psi \rangle - 1)}{\int d[\psi] \delta(\langle \chi | \psi \rangle - z) \delta(\langle \psi | \psi \rangle - 1)}$$

Recipe for integration

$$\text{HA}_\psi \left[|\langle \chi | A | \psi \rangle|^2 \delta(\langle \chi | \psi \rangle - z) \right] = \frac{\int d[\psi] |\langle \chi | A | \psi \rangle|^2 \delta(\langle \chi | \psi \rangle - z) \delta(\langle \psi | \psi \rangle - 1)}{\int d[\psi] \delta(\langle \chi | \psi \rangle - z) \delta(\langle \psi | \psi \rangle - 1)}$$

Generating function and Fourier integrals

$$|\langle \chi | A | \psi \rangle|^2 = \frac{\partial_J}{i} \text{Exp } iJ \langle \chi | A | \psi \rangle \langle \psi | A^\dagger | \chi \rangle \Big|_{J=0}$$
$$\delta(x) = \frac{1}{2\pi} \int_{\mathbb{R}} dt \text{ Exp}[itx]$$

⇒ Arrive at Gaussian integral

Transition probabilities

For the Hilbert space average- and variance we find

$$\text{HA}_\psi [|\langle \chi | A | \psi \rangle|^2 \delta(\langle \chi | \psi \rangle - z)] = \frac{1 - |z|^2}{N - 1} \langle \chi | A A^\dagger | \chi \rangle + \frac{N |z|^2 - 1}{N - 1} |\langle \chi | A | \chi \rangle|^2$$

$$\begin{aligned} \text{HV}_\psi \left[|\langle \chi | A | \psi \rangle|^2 \delta(\langle \chi | \psi \rangle - z) \right] &= \lambda_1(N, |z|) \langle \chi | A A^\dagger | \chi \rangle^2 + \lambda_2(N, |z|) \langle \chi | A A^\dagger | \chi \rangle \\ &\quad + \lambda_3(N, |z|) |\langle \chi | A | \chi \rangle|^4 \end{aligned}$$

Special cases

$$\text{HA}_\psi [|\langle \chi | A | \psi \rangle|^2 \delta(\langle \chi | \psi \rangle - 1)] = |\langle \chi | A | \chi \rangle|^2$$

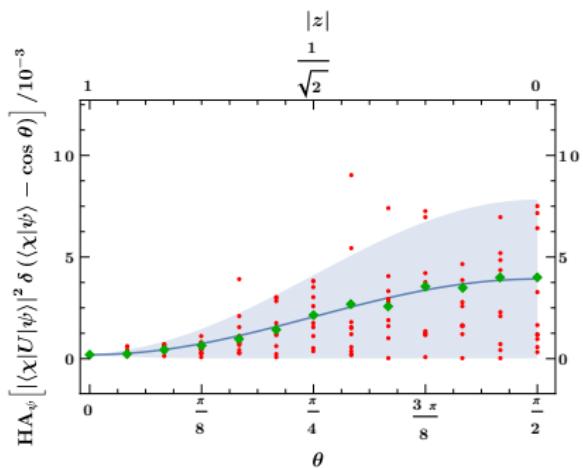
$$\text{HV}_\psi [|\langle \chi | A | \psi \rangle|^2 \delta(\langle \chi | \psi \rangle - 1)] = 0$$

$$\text{HA}_\psi [|\langle \chi | \psi \rangle|^2 \delta(\langle \chi | \psi \rangle - z)] = |z|^2$$

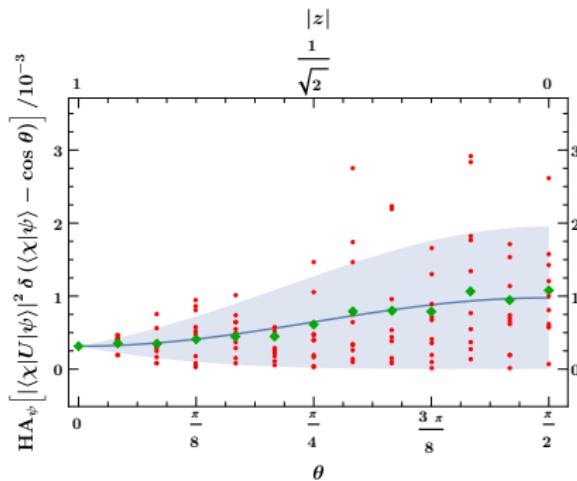
$$\text{HV}_\psi [|\langle \chi | \psi \rangle|^2 \delta(\langle \chi | \psi \rangle - z)] = 0$$

Transition probabilities

Let $A = U$



$n = 8$

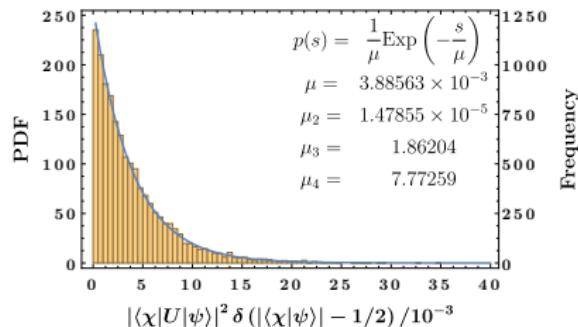
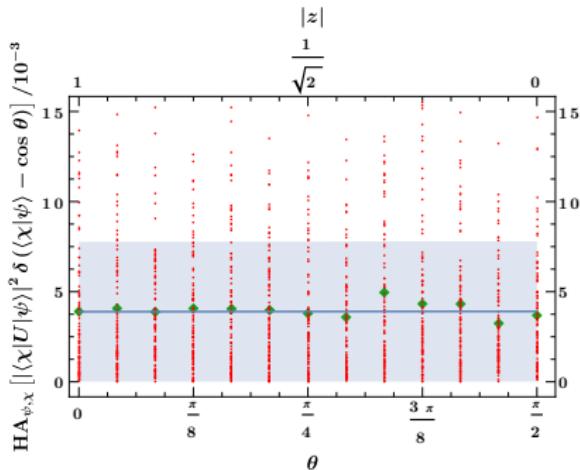


$n = 10$

$$\text{HA}_\psi [|\langle \chi | U | \psi \rangle|^2 \delta(\langle \chi | \psi \rangle - z)] = \frac{1 - |z|^2}{N - 1} + \frac{N |z|^2 - 1}{N - 1} |\langle \chi | U | \chi \rangle|^2$$

Transition probabilities

Now average over the remaining state $|\chi\rangle$



$$\text{HA}_{\psi,\chi} \left[|\langle \chi | U | \psi \rangle|^2 \delta(\langle \chi | \psi \rangle - z) \right] = \frac{N - |z|^2}{N^2 - 1} + \frac{N |z|^2 - 1}{N^3 - N} |\text{Tr } U|^2$$

$$\text{HV}_{\psi,\chi} \left[|\langle \chi | U | \psi \rangle|^2 \delta(\langle \chi | \psi \rangle - z) \right] \approx \text{const.}$$

Spectral Form Factor

With Spectral Form Factor $K(T) = \frac{|\text{Tr } U^T|^2}{N}$

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Dual Operator W with $\dim W = 2^T \times 2^T$

$$\text{Tr } U^T = \text{Tr } W^n$$

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Dual Operator W with $\dim W = 2^T \times 2^T$

$$\text{Tr } U^T = \text{Tr } W^n$$

If n large replace $\text{Tr } W^n$ by the largest eigenvalue $\tilde{\lambda}_{\max}^n(J, b, h)$

$$\text{HA}_{\psi,\chi} \left[|\langle \chi | U | \psi \rangle|^2 \delta(\langle \chi | \psi \rangle - z) \right] = \frac{N - |z|^2}{N^2 - 1} + \frac{N|z|^2 - 1}{N^2 - 1} \left| \tilde{\lambda}_{\max}^n \right|^2$$

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Connection to spectral statistics

What happens for non-uniformly distributed $|\psi\rangle$ and $|\chi\rangle$?

Thank you for the attention!